ELECTRONIC JOURNAL OF DIFFERENTIAL EQUATIONS, Vol. **1998**(1998), No. 16, pp. 1–10. ISSN: 1072-6691. URL: http://ejde.math.swt.edu or http://ejde.math.unt.edu ftp (login: ftp) 147.26.103.110 or 129.120.3.113

Quasi-geostrophic type equations with weak initial data *

Jiahong Wu

Abstract

We study the initial value problem for the quasi-geostrophic type equations

$$\begin{aligned} \frac{\partial \theta}{\partial t} + u \cdot \nabla \theta + (-\Delta)^{\lambda} \theta &= 0, \quad \text{on} \quad \mathbb{R}^{n} \times (0, \infty), \\ \theta(x, 0) &= \theta_{0}(x), \quad x \in \mathbb{R}^{n}, \end{aligned}$$

where $\lambda(0 \leq \lambda \leq 1)$ is a fixed parameter and $u = (u_j)$ is divergence free and determined from θ through the Riesz transform $u_j = \pm \mathcal{R}_{\pi(j)}\theta$, with $\pi(j)$ a permutation of $1, 2, \dots, n$. The initial data θ_0 is taken in the Sobolev space $\dot{L}_{r,p}$ with negative indices. We prove local well-posedness when

$$\frac{1}{2} < \lambda \leq 1, \quad 1 < p < \infty, \quad \frac{n}{p} \leq 2\lambda - 1, \quad r = \frac{n}{p} - (2\lambda - 1) \leq 0.$$

We also prove that the solution is global if θ_0 is sufficiently small.

1 Introduction

In this paper we study the initial value problem (IVP) of the dissipative quasigeostrophic type (QGS) equations

$$\frac{\partial\theta}{\partial t} + u \cdot \nabla\theta + (-\Delta)^{\lambda}\theta = 0, \quad \text{on} \quad \mathbb{R}^n \times (0, \infty), \tag{1.1}$$

$$\theta(x,0) = \theta_0(x), \quad x \in \mathbb{R}^n \tag{1.2}$$

where $\lambda(0 \leq \lambda \leq 1)$ is a fixed parameter and the velocity $u = (u_1, u_2, \dots, u_n)$ is divergence free and determined from θ by

$$u_j = \pm \mathcal{R}_{\pi(j)} \theta, \qquad \pi(j) \text{ is a permutation of } 1, 2, \cdots, n$$
 (1.3)

^{*1991} Mathematics Subject Classifications: 35K22, 35Q35, 76U05.

Key words and phrases: Quasi-geostrophic equations, Weak data, Well-posedness.

^{©1998} Southwest Texas State University and University of North Texas.

Submitted November 26, 1996. Published June 12, 1998.

Supported by NSF grant DMS 9304580 at IAS.

where u_j may take either + or - sign and $\mathcal{R}_j = \partial_j (-\Delta)^{-1/2}$ are the Riesz transforms. Here Riesz potential operator $(-\Delta)^{\alpha}$ is defined through the Fourier transform:

$$\widehat{f}(\xi) = \int e^{-2\pi i x \cdot \xi} f(x) \, dx$$
$$(\widehat{(-\Delta)^{\alpha}} f)(\xi) = (2\pi |\xi|)^{2\alpha} \widehat{f}(\xi)$$

A particularly important special case of (1.1) is the 2-D dissipative quasigeostrophic equations in which the velocity $u = (u_1, u_2)$ can also be defined through the stream function ψ :

$$u = (u_1, u_2) = \left(-\frac{\partial \psi}{\partial x_2}, \frac{\partial \psi}{\partial x_1}\right), \quad (-\Delta)^{1/2} \psi = -\theta \tag{1.4}$$

The 2-D QGS equations are derived from more general quasi-geostrophic approximations for flow in rapidly rotating 3-D half space, which in some important cases reduce to the evolution equation for the temperature on the 2-D boundary given in (1.1), (1.2),(1.4) ([12, 2]). The scalar θ represents the potential temperature and u is the fluid velocity. These equations have been under active investigation because of mathematical importance and potential applications in meteorology and oceanography ([12, 2, 1, 6]). As pointed out in [2], the non-dissipative 2-D QGS equations are strikingly analogous to the 3-D Euler equations and thus serve as a simple model in seeking possible singular solutions.

We are interested mainly in the well-posedness result for initial data θ_0 in homogeneous Lebesgue spaces, $\theta_0 \in \dot{L}_{r,p}(\mathbb{R}^n)$ (defined below). By wellposedness we mean existence, uniqueness and persistence (i.e. the solution describes a continuous curve belonging to the same space as does the initial data) and continuous dependence on the data.

Here the homogeneous Lebesgue space $\dot{L}_{s,q}(\mathbb{R}^n)$ consists of all v such that

$$(-\Delta)^{\frac{s}{2}} v \in L^q, \quad s \in \mathbb{R}, \quad 1 \le q < \infty,$$

and the standard norm is given by

$$||v||_{s,q} = ||(-\Delta)^{s/2}v||_{L^q}.$$

These spaces are also called the spaces of Riesz potentials. Kato and Ponce [10] consider the Navier-Stokes equations with initial data in this type of spaces.

We prove that if $\frac{1}{2} < \lambda \leq 1$ and $\theta_0 \in L_{r,p}$ with r, p satisfying

$$1$$

then the IVP (1.1), (1.3), (1.2) is locally well-posed. The solution is global if θ_0 is sufficiently small. The detailed statements are given in Theorem 2.2 of the next section.

Jiahong Wu

Although there is a large body of literature on quasi-geostrophic equations ([12, 1, 6, 2]), not many rigorous mathematical results concerning the solutions have been obtained. In [2] Constantin-Majda-Tabak proved finite time existence results for smooth data and developed mathematical criteria characterizing blowup for the 2-D non-dissipative QGS equation. In [13] Resnick obtained solutions of 2-D QGS equations with L^2 data on periodic domain by using Galerkin approximation. In a previous paper [15], the vanishing dissipation limits and Gevrey class regularity [3] for the 2-D dissipative QGS equations are obtained. In this paper we consider the IVP of the general *n*-D QGS type equations (defined by (1.1), (1.3, (1.2)) with initial data in Sobolev spaces of negative indices and establish local well-posedness results. For sufficiently small initial data, the solution is global. By taking n = 2 and p = 2, the well-posedness reduces to the L^2 results in 2-D.

The main result is presented in the next section, and it is proven using the contraction-mapping principle.

2 Well-posedness

We need to use the spaces of weighted continuous functions in time, which have been introduced by Kato, Ponce and others in solving the Navier-Stokes equations ([8, 10, 11]).

Definition 2.1 Suppose T > 0 and $\alpha \ge 0$ are real numbers. The spaces $C_{\alpha,s,q}$ and $\dot{C}_{\alpha,s,q}$ are defined as

$$C_{\alpha,s,q} \equiv \{ f \in C((0,T), L_{s,q}), \| f \|_{\alpha,s,q} < \infty \},\$$

where the norm is given by

$$||f||_{\alpha,s,q} = \sup\{t^{\alpha}||f||_{s,q}, \quad t \in (0,T)\}.$$

Note that $\dot{C}_{\alpha,s,q}$ is a subspace of $C_{\alpha,s,q}$:

$$\dot{C}_{\alpha,s,q} \equiv \left\{f \in C_{\alpha,s,q}, \quad \lim_{t \to 0} t^\alpha \|f(t)\|_{s,q} = 0\right\}.$$

When $\alpha = 0$, the spaces $\bar{C}_{s,q}$ are used for $BC([0,T), \dot{L}_{s,q})$.

These spaces are important in uniqueness and local existence problems ([8, 10, 11]). Notice that $f \in C_{\alpha,s,q}$ (resp. $f \in \dot{C}_{\alpha,s,q}$) implies that $||f(t)||_{s,q} = O(t^{-\alpha})$ (resp. $o(t^{-\alpha})$).

The main result of this section is the well-posedness theorem that states

Theorem 2.2 Assume that $\lambda > 1/2$ and $\theta_0 \in L_{r,p}$ with r, p satisfying

$$1 (2.1)$$

Then there exists $T = T(\theta_0)$ and a unique solution $\theta(t)$ of the IVP (1.1),(1.3), (1.2) in the time interval [0,T) satisfying

$$\theta \in Y_T \equiv \left(\bigcap_{p \le q < \infty} \bar{C}_{\frac{n}{q} - (2\lambda - 1), q}\right) \cap \left(\bigcap_{p \le q < \infty} \bigcap_{s > \frac{n}{q} - (2\lambda - 1)} \dot{C}_{\left(s - \frac{n}{q} + (2\lambda - 1)\right)/(2\lambda), s, q}\right)$$

In particular,

$$\theta \in BC([0,T), \dot{L}_{r,p}) \cap (\cap_{s>r}C((0,T), \dot{L}_{s,p}))$$

Furthermore, for some neighborhood V of θ_0 , the mapping

$$\mathfrak{P}: V \longmapsto Y_T: \quad \theta_0 \longmapsto \theta$$

is Lipschitz.

Remark 2.3 If $\|\theta_0\|_{r,p}$ is small enough, then we can take $T = \infty$.

We prove this theorem by the method of integral equations and contractionmapping arguments. Following standard practice ([4, 5, 7, 10]), we write the QGS equation (1.1) into the integral form:

$$\theta = K\theta_0(t) - G(u,\theta)(t) \equiv e^{-\Lambda^{2\lambda}t}\theta_0 - \int_0^t e^{-\Lambda^{2\lambda}(t-\tau)} (u \cdot \nabla\theta)(\tau) d\tau , \qquad (2.2)$$

where $K(t) = e^{-\Lambda^{2\lambda}t}$ is the solution operator of the linear equation

$$\partial_t \theta + \Lambda^{2\lambda} \theta = 0$$
, with $\Lambda = (-\Delta)^{1/2}$

We observe that $u \cdot \nabla \theta = \sum_{j} u_{j} \partial_{j} \theta = \nabla \cdot (u\theta)$ provided that $\nabla \cdot u = 0$. This provides an alternative expression for G:

$$G(u,\theta)(t) = G(u\theta)(t) = \int_0^t \nabla \cdot e^{-\Lambda^{2\lambda}(t-\tau)}(u\theta)(\tau)d\tau$$

We shall solve (2.2) in the spaces of weighted continuous functions in time introduced in the beginning of this section. To this end we need estimates for the operators K and G acting between these spaces. These are established in the two propositions that follow.

Proposition 2.4 (i) For $1 \le q < \infty$ and $s \in \mathbb{R}$, the operator K maps continuously from $\dot{L}_{s,q}$ into $\bar{C}_{s,q} \equiv BC([0,\infty), \dot{L}_{s,q})$.

(ii) If q_1, q_2, s_1, s_2 and α_2 satisfy $q_1 \leq q_2$, $s_1 \leq s_2$, and

$$\alpha_2 = \frac{1}{2\lambda}(s_2 - s_1) + \frac{1}{2\lambda}\left(\frac{n}{q_1} - \frac{n}{q_2}\right),$$

then K maps continuously from \dot{L}_{s_1,q_1} to $\dot{C}_{\alpha_2,s_2,q_2}$ (When $\alpha_2 = 0$, \dot{C} should be replaced by \bar{C}).

EJDE-1998/16

Jiahong Wu

Proof. To prove Assertion (i), it suffices to prove that for some constant C,

$$\|K\phi(t)\|_{L^q} \le C \|\phi\|_{L^q}, \quad \text{for any } t \in [0,\infty),$$

which can be established using the Young's inequality

$$||K\phi(t)||_{L^q} \le ||K(t)||_{L^1} ||\phi||_{L^q}$$

and the fact that

with $\frac{1}{q} = 1$

$$\widehat{K}(t)(\xi) = e^{-|2\pi\xi|^{2\lambda}t}, \quad ||K(t)||_{L^1} = \widehat{K}(t)(0) = 1.$$

To prove Assertion (ii), we first note that the operator $(-\Delta)^{s_0/2}K(t)$ has the property

$$\|(-\Delta)^{s_0/2} K(t)\|_{L^q(\mathbb{R}^n)} \le C t^{\frac{1}{2\lambda} \left(-s_0 - n(1 - \frac{1}{q})\right)}, \qquad (2.3)$$

where $s_0 \ge 0$, $q \in [1, \infty)$ and C is a constant. The proof of this property is similar to that for the heat operator ([4, 5, 10]). To show (ii), it suffices show that for some constant C,

$$\sup_{t \in [0,T)} t^{\alpha_2} \| (-\Delta)^{\frac{s_0}{2}} K \phi(t) \|_{L^{q_2}} \le C \| \phi \|_{L^{q_1}}$$

with $s_0 = s_2 - s_1 \ge 0$. This can be proved using the property (2.3) and Young's inequality

$$\|(-\Delta)^{\frac{s_0}{2}} K\phi(t)\|_{L^{q_2}} \le C \|(-\Delta)^{\frac{s_0}{2}} K(t)\|_{L^q} \|\phi\|_{L^{q_1}} - \left(\frac{1}{q_1} - \frac{1}{q_2}\right). \qquad \Box$$

Now we give estimates for the operator

$$G(g)(t) = \int_0^t \nabla \cdot K(t-\tau)g(\tau)d\tau$$

Proposition 2.5 If $q_1, q_2, s_1, s_2, \alpha_1$ and α_2 satisfy $q_1 \leq q_2$,

$$\begin{split} s_1 - 1 &\leq s_2 < s_1 + 2\lambda - 1 - \left(\frac{n}{q_1} - \frac{n}{q_2}\right) \\ \alpha_1 < 1, \quad and \quad \alpha_2 &= \alpha_1 - 1 + \frac{1}{2\lambda} \left[s_2 - s_1 + 1 + \frac{n}{q_1} - \frac{n}{q_2}\right] \,, \end{split}$$

then G is a continuous mapping from $\dot{C}_{\alpha_1,s_1,q_1}$ to $\dot{C}_{\alpha_2,s_2,q_2}$.

Proof. Let $g \in \dot{C}_{\alpha_1,s_1,q_1}$. Then clearly,

$$\|G(g)\|_{\alpha_{2},s_{2},q_{2}} = \sup_{t \in [0,T)} t^{\alpha_{2}} \int_{0}^{t} \|(-\Delta)^{\frac{(1+s_{0})}{2}} K(t-\tau) \left((-\Delta)^{\frac{s_{1}}{2}} g(\tau)\right)\|_{L^{q_{2}}} d\tau$$

where $s_0 = s_2 - s_1$. Using Young's inequality,

$$\|G(g)\|_{\alpha_{2},s_{2},q_{2}} \leq \sup_{t \in [0,T)} t^{\alpha_{2}} \int_{0}^{t} \|(-\Delta)^{\frac{(1+s_{0})}{2}} K(t-\tau)\|_{L^{q}} \|\left((-\Delta)^{\frac{s_{1}}{2}} g(\tau)\right)\|_{L^{q_{1}}} d\tau$$

with $\frac{1}{q} = 1 - \left(\frac{1}{q_1} - \frac{1}{q_2}\right)$. If $s_0 + 1 \ge 0$, we can use the property (2.3) of operator K and obtain

$$\begin{split} \|G(g)\|_{\alpha_{2},s_{2},q_{2}} &\leq C \|g\|_{\alpha_{1},s_{1},q_{1}} \sup_{t \in [0,T)} t^{\alpha_{2}} \int_{0}^{t} (t-\tau)^{-\frac{1}{2\lambda} \left(s_{0}+1+n(1-\frac{1}{q})\right)} \tau^{-\alpha_{1}} d\tau \\ &\leq C \|g\|_{\alpha_{1},s_{1},q_{1}} \sup_{t \in [0,T)} t^{\alpha_{2}-\alpha_{1}+1-\frac{1}{2\lambda} \left(s_{0}+1+n(1-\frac{1}{q})\right)} \times \\ & B\left(1-\frac{1}{2\lambda} \left[s_{0}+1+n(1-\frac{1}{q})\right], 1-\alpha_{1}\right), \end{split}$$

where C is a constant and B(a, b) is the Beta function

$$B(a,b) = \int_0^1 (1-x)^{a-1} x^{b-1} \, dx \, .$$

By noticing that B(a, b) is finite when a > 0, b > 0 and that

$$s_0 = s_2 - s_1, \quad 1 - \frac{1}{q} = \frac{1}{q_1} - \frac{1}{q_2}$$

we obtain

$$\|G(g)\|_{\alpha_2, s_2, q_2} \le C \|g\|_{\alpha_1, s_1, q_1},$$

if the indices satisfy $0 \le s_2 - s_1 + 1 < 2\lambda - \frac{n}{q_1} - \frac{n}{q_2}$, $\alpha_1 < 1$, and $\alpha_2 = \alpha_1 - 1 + \frac{1}{2\lambda} \left[s_2 - s_1 + 1 + \frac{n}{q_1} - \frac{n}{q_2} \right]$.

$$-1 + rac{1}{2\lambda} \left[s_2 - s_1 + 1 + rac{n}{q_1} - rac{n}{q_2}
ight].$$

To prove Theorem 2.2, we also need the following singular integral operator estimate whose proof can be found in [14].

Lemma 2.6 For $u = (u_j)$ with $u_j = \pm \mathcal{R}_{\pi(j)}\theta(j = 1, 2, \dots, n)$, where \mathcal{R}_j are the Riesz transforms, we have the estimate

$$\|u\|_{L^q} \le C_q \|\theta\|_{L^q}, \quad 1 < q < \infty$$

with C_q a constant depending on q.

Proof of Theorem 2.2. We distinguish between two cases: r < 0, and r = 0. For r < 0, we define

$$X = \bar{C}_{r,p} \cap \dot{C}_{-\frac{r}{2\lambda},0,p}$$

with norm for $\theta \in X$ given by

$$\|\theta\|_X = \|\theta - K\theta_0\|_{0,r,p} + \|\theta\|_{-\frac{r}{2\lambda},0,p}$$

and the complete metric space X_R to be the closed ball in X of radius R. Consider the operator $\mathcal{A}(\theta, \theta_0) : X_R \times V \longmapsto X$

$$\mathcal{A}(\theta, \theta_0)(t) = K\theta_0(t) - G(u\theta)(t), \quad 0 < t < T,$$

where V is some neighborhood of θ_0 in $\dot{L}_{r,p}$ and T will be chosen. Using Proposition 2.4 by substituting s = r, q = p in (i) and

$$q_1 = q_2 = p, \quad s_1 = r, \quad s_2 = 0, \quad \alpha_2 = -\frac{r}{2\lambda}$$

in (ii), we find that $K\tilde{\theta}_0(t) \in X_R$ for $\tilde{\theta}_0 \in V$ if T is taken small enough and V is chosen properly.

To estimate G, we use Proposition 2.5 with

$$q_1 = \frac{p}{2}, \quad q_2 = p, \quad s_1 = 0, \quad s_2 = l + r, \quad \alpha_1 = -\frac{r}{\lambda}, \quad \alpha_2 = \frac{l}{2\lambda}$$

to obtain for a constant c such that

$$\|G(u\theta)\|_{\frac{l}{2\lambda},l+r,p} \le c \|u\theta\|_{-\frac{r}{\lambda},0,\frac{p}{2}} \le c \|u\|_{-\frac{r}{2\lambda},0,p} \|\theta\|_{-\frac{r}{2\lambda},0,p}$$

for $l \in [0, -2r)$. To estimate u in terms of θ , we use Lemma 2.6, i.e. for 1 ,

$$\|u\|_{L^p} \le C_p \|\theta\|_{L^p}$$

and eventually we obtain

$$\|G(u\theta)\|_{\frac{l}{2\lambda}, l+r, p} \le cC_p \|\theta\|_{-\frac{r}{2\lambda}, 0, p}^2 \le cC_p R^2.$$

Notice that the restrictions (2.1) on r, p are necessary in order to apply Propositions 2.4, 2.5 and Lemma 2.6.

Furthermore,

$$\|\mathcal{A}(\theta,\theta_0) - \mathcal{A}(\tilde{\theta},\theta_0)\|_X = \|G(u\theta) - G(\tilde{u}\tilde{\theta})\|_X,$$

~

where $\tilde{u} = (\tilde{u}_j)$ with $\tilde{u}_j = \pm \mathcal{R}_{\pi(j)} \tilde{\theta}(j = 1, 2, \cdots, n)$. Using Proposition 2.5 again,

$$\begin{aligned} \|\mathcal{A}(\theta,\theta_0) - \mathcal{A}(\tilde{\theta},\theta_0)\|_X &\leq \|G((\tilde{u}-u)\tilde{\theta})\|_X + \|G(u(\theta-\tilde{\theta}))\|_X \\ &\leq c\left(\|\tilde{u}-u\|_X\|\tilde{\theta}\|_X + \|\theta-\tilde{\theta}\|_X\|u\|_X\right). \end{aligned}$$

Since $(\tilde{u} - u)_j = \pm \mathcal{R}_{\pi(j)}(\tilde{\theta} - \theta)$, Lemma 2.6 implies

$$\|u\|_X \le C_p \|\theta\|_X, \quad \|\tilde{u} - u\|_X \le C_p \|\tilde{\theta} - \theta\|_X.$$

Therefore, for constant satisfies $C = cC_p$ and

$$\|\mathcal{A}(\theta,\theta_0) - \mathcal{A}(\theta,\theta_0)\|_X \le C(\|\theta\|_X + \|\theta\|_X)\|\theta - \theta\|_X.$$

Our above estimates show that if we choose T small and R appropriately, then \mathcal{A} maps X_R into itself and is a contraction. Consequently there exists a unique fixed point $\theta \in X_R$: $\theta = \mathfrak{P}(\theta_0)$ satisfying $\theta = \mathcal{A}(\theta, \theta_0)$. It is easy to see from these estimates that the uniqueness can be extended to all R' by further reducing the the time interval and thus to the whole X.

To prove the Lipschitz continuity of \mathfrak{P} on V, let $\theta = \mathfrak{P}(\theta_0)$ and $\zeta = \mathfrak{P}(\zeta_0)$ for $\theta_0, \zeta_0 \in V$. Then

$$\begin{split} \|\theta - \zeta\|_X &= \|\mathcal{A}(\theta, \theta_0) - \mathcal{A}(\zeta, \zeta_0)\|_X \\ &\leq \quad \|\mathcal{A}(\theta, \theta_0) - \mathcal{A}(\zeta, \theta_0)\|_X + \|\mathcal{A}(\zeta, \theta_0) - \mathcal{A}(\zeta, \zeta_0)\|_X \\ &\leq \quad \gamma \|\theta - \zeta\|_X + \|K(\theta_0 - \zeta_0)\|_X \end{split}$$

Since \mathcal{A} is a contraction, $\gamma < 1$. Therefore, the asserted property is obtained by applying Proposition 2.4 to the second term of the last inequality.

To show that θ is in the asserted class Y_T (defined in Theorem 2.2), we notice that

$$\theta = \mathcal{A}(\theta, \theta_0) \equiv K\theta_0 - G(u\theta).$$

We apply Proposition 2.4 twice to $K\theta_0$ to show that

$$K\theta_0 \in \bar{C}_{\frac{n}{q}-(2\lambda-1),q}, \qquad K\theta_0 \in \dot{C}_{\left(s-\frac{n}{q}+(2\lambda-1)\right)/(2\lambda),s,q}$$

for any $p \leq q < \infty$ and $s > \frac{n}{q} - (2\lambda - 1)$. To show the second part

$$G(u\theta) \in \bar{C}_{\frac{n}{q}-(2\lambda-1),q}, \quad p \le q < \infty$$
(2.4)

we use Proposition 2.5 with

$$q_1 = \frac{p}{2}, \quad q_2 = q, \quad s_1 = 0, \quad s_2 = \frac{n}{q} - (2\lambda - 1), \quad \alpha_1 = -\frac{r}{\lambda}, \quad \alpha_2 = 0$$

and obtain

$$\|G(u\theta)\|_{0,\frac{n}{q}-(2\lambda-1),q} \le C \|u\theta\|_{-\frac{r}{\lambda},0,\frac{p}{2}} \le C \|u\|_{-\frac{r}{2\lambda},0,p} \|\theta\|_{-\frac{r}{2\lambda},0,p}.$$

The asserted property (2.4) is established after we apply Lemma 2.6 to u. Once again, we apply Proposition 2.5 with

$$q_1 = \frac{p}{2}, \quad q_2 = q, \quad s_1 = 0, \quad s_2 = s,$$
$$\alpha_1 = -\frac{r}{\lambda}, \quad \alpha_2 = \frac{1}{2\lambda} \left[s - \left(\frac{n}{q} - (2\lambda - 1)\right) \right]$$

8

EJDE-1998/16

Jiahong Wu

to show that

$$G(u\theta) \in \dot{C}_{\left(s-\frac{n}{q}+(2\lambda-1)\right)/(2\lambda),s,q}, \quad \text{for } s > \frac{n}{q} - (2\lambda-1), \tag{2.5}$$

but s should also satisfy

$$s < 2\lambda - 1 - \left(\frac{2n}{p} - \frac{n}{q}\right)$$

as required by Proposition 2.5. For large s, (2.5) can be shown by an induction process (see an analogous argument in [8]).

We now deal with the case r = 0. Define

$$X = \bar{C}_{0,p} \cap \dot{C}_{\frac{1}{4},0,\frac{4\lambda-2}{3\lambda-2}p}$$

with the norm

$$\|\theta\|_X = \|\theta - K\theta_0\|_{0,0,p} + \|\theta\|_{\frac{1}{4},0,\frac{4\lambda-2}{3\lambda-2}p}.$$

For $\theta \in X_R$, we have by Proposition 2.5,

$$\begin{aligned} \|G(u\theta)\|_{X} &= \|G(u\theta)\|_{0,0,p} + \|G(u\theta)\|_{\frac{1}{4},0,\frac{4\lambda-2}{3\lambda-2}p} \\ &\leq c \|u\theta\|_{\frac{1}{2},0,\frac{2\lambda-1}{3\lambda-2}p} \\ &\leq c \|u\|_{\frac{1}{4},0,\frac{4\lambda-2}{3\lambda-2}p} \|\theta\|_{\frac{1}{4},0,\frac{4\lambda-2}{3\lambda-2}p} \,. \end{aligned}$$

Here c is a constant which may depend on the indices λ , p, and n. Using Lemma 2.6 again, we obtain a constant C such that

$$\|G(u\theta)\|_X \le C \|\theta\|_X^2 \le CR^2.$$

Once the above estimates have been established, the rest of the proof in this case is similar to that described in the case r < 0.

Acknowledgments I would like to thank Professor P. Constantin for teaching me the quasi-geostrophic equations and Professor C. Kenig for his helpful suggestions.

References

- J. Charney, N. Phillips, Numerical integrations of the quasi-geostrophic equations for barotropic and simple baroclinic flows, J. Meteorol., 10(1953), 71-99.
- [2] P. Constantin, A. Majda, E. Tabak, Formation of strong fronts in the 2-D quasi-geostrophic thermal active scalar, Nonlinearity, 7(1994), 1495-1533

- [3] C. Foias, R. Temam, Gevrey class regularity for the solutions. of the Navier-Stokes equations, J. Funct. Anal., 87(1989), 359-369.
- [4] Y. Giga, Solutions for semilinear parabolic equations in L^p and regularity of weak solutions of the Navier-Stokes system, J. Diff. Eq., 62(1986), 186-212.
- [5] Y. Giga, T. Miyakawa, H. Osada, Two dimensional Navier-Stokes flow with measures as initial vorticity, Arch. Rational Mech. Anal., 104(1988), 223-250.
- [6] I. Held, R. Pierrehumbert, S. Garner, K. Swanson, Surface quasigeostrophic dynamics, J. Fluid Mech., 282(1995), 1-20.
- [7] T. Kato, Strong L^p solutions of the Navier-Stokes equation in \mathbb{R}^m with applications to weak solutions, Math. Z., **187**(1984), 471-480.
- [8] T. Kato, The Navier-Stokes solutions for an incompressible fluid in ℝ² with measure as the initial vorticity, Diff. Integral Eq., 7(1994), 949-966.
- [9] T. Kato, H. Fujita, On the nonstationary Navier-Stokes equations, Rend. Sem. mat. Univ. Padova, 32(1962), 243-260.
- [10] T. Kato, G. Ponce, The Navier-Stokes equations with weak initial data, Int. Math. Research Notices, 10(1994), 435-444.
- [11] T. Kato, G. Ponce, Well-posedness of the Euler and the Navier-Stokes equations in the Lebesgue spaces L^p_s(R²), Rev. Mat. Iber. 2(1986), 73-88.
- [12] J. Pedlosky, "Geophysical Fluid Dynamics", Springer, New York, 1987.
- [13] S. Resnick, "Dynamical Problems in Non-linear Advective Partial Differential Equations", Ph.D. thesis, University of Chicago, 1995
- [14] E.M. Stein, "Singular Integrals and Differentiability Properties of Functions", NJ: Princeton University Press, 1970
- [15] J. Wu, Inviscid limits and regularity estimates for the solutions of the 2-D dissipative quasi-geostrophic equations, Indiana Univ. Math. J., 46(1997), in press.

JIAHONG WU School of Mathematics, Institute for Advanced Study Princeton, NJ 08540. USA E-mail address: jiahong@math.utexas.edu