# Continuous Piecewise Polynomials and Static Equilibrium 

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Partition $\mathcal{P}$ of an octagonal domain $\subset \mathbb{R}^{2}$
Graph of the ZwartPowell element: a spline in $C_{2}^{1}(\mathcal{P})$

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Graph of PL function on I

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Trapezoid Rule!

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Simpson's Rule!

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- ... the list goes on


## Univariate Piecewise Linear Functions

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Basis for $C_{1}^{0}(I)$ : 'Courant functions' or 'tent functions' which are 1 at a chosen vertex and 0 at all others.

## Tent Functions 1

Univariate Courant functions:


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Can generalize this dimension formula for all $r, d$ :

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\operatorname{dim}_{\mathbb{R}} C_{d}^{r}(I)=\left\{\begin{array}{rl}
d+1 & d \leq r \\
e(d+1)-v^{0}(r+1) & d>r
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There are nice algorithms due to Casteljau and de Boor to compute bases of $C_{d}^{r}(I)$ called $\mathbf{B}$-splines.

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A triangulation $\Delta$ with $v=8, e=15, f=8, v^{0}=2$, and $e^{0}=9$
What kinds of PL functions are there on $\Delta$ ?

Again, a continuous piecewise linear function on $\Delta$ is uniquely determined by its value on the vertices (3 points determine a plane!).


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Just as in the 1D case, $C_{1}^{0}(\Delta)$ is a vector space. Again, a basis for $C_{1}^{0}(\Delta)$ is given by the 'Courant functions' which are 1 at a chosen vertex and 0 on all other vertices.

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Plugging in $d=0$ gives 1 , plugging in $d=1$ gives $v$. (this takes a little work)
There is no reference to the geometry of $\Delta$ ! All that matters is the number of faces, edges, and vertices.

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- Call this a polygonal graph or polygonal framework
- $f, e, v, e^{0}, v^{0}$ stay the same


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A polygonal framework $\mathcal{P}_{1}$ with

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Let's see why.

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When you move to $\mathcal{P}_{2}$ you lose this PL function!

## Generating Interesting Examples

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Make it transparent Now look in one of the faces:


The nontrivial PL function is a 'deformed cube'

## A more interesting example

Chop off cube corners


## A more interesting example

Make it transparent


## A more interesting example

 Make it transparent Look into an octagonal face:

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Make it transparent Look into an octagonal face:


We get a nontrivial PL function which is a 'deformed' version of the truncated cube

## And now for something completely different

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Between vertices $p_{i}, p_{j}$ (thought of as vectors), represent tension or compression as a scalar $\omega_{i j}$.

- Force is $\omega_{i j}\left(p_{j}-p_{i}\right)$ at $p_{i}$
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- Force is $\omega_{i j}\left(p_{j}-p_{i}\right)$ at $p_{i} \vee \omega_{i j}<0 \Longrightarrow$ tension
- Force is $\omega_{i j}\left(p_{i}-p_{j}\right)$ at $p_{j} \vee \omega_{i j}>0 \Longrightarrow$ compression


## Self-Stress

A self-stress on a framework is an assignment of scalars $\omega_{i j}$ along the edges $e_{i j}$ satisfying

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By the way, what could this mean physically?

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Fact: If the domain is not simply connected, the above correspondence breaks down!

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$\widehat{\mathcal{P}}_{1}$
- $C^{0}(\widehat{\mathcal{P}})$ is graded (every spline can be written as a sum of splines of uniform degree)
- $C_{d}^{0}(\mathcal{P})$ 'sits inside' $C^{r}(\widehat{\mathcal{P}})$ as the degree $d$ 'slice.'


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- Useful to consider algebraic structures on $C^{0}(\widehat{\mathcal{P}})$ in addition to vector space structure
- $F \in C^{0}(\widehat{\mathcal{P}}), f \in \mathbb{R}[x, y, z]$ a polynomial. Then $f \cdot F \in C^{0}(\widehat{\mathcal{P}})$.
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- Relation between PL functions and self-stresses generalizes to a correspondence between splines and syzygies, and dependence of this correspondence on $\mathcal{P}$ being simply connected is completely clarified (thanks to Billera).
- Via some homological algebra, $\operatorname{dim} C_{1}^{0}(\mathcal{P})$ has consequences for freeness of $C^{0}(\widehat{\mathcal{P}})$ as an $\mathbb{R}[x, y, z]$-module. This in turn impacts how easy it is to calculate $\operatorname{dim} C_{d}^{0}(\mathcal{P})$ for $d \geq 1$.

THANK YOU!

